

1 Minimum relative entropy, Bayes, and Kapur

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3 4 SUMMARY

5
6 The focus of this paper is to illustrate important philosophies on inversion and the
7 similarly and differences between Bayesian and minimum relative entropy (MRE)
8 methods. The development of each approach is illustrated through the general-
9 discrete linear inverse. MRE differs from both Bayes and classical statistical meth-
10 ods in that knowledge of moments are used as “data” rather than sample values.
11 MRE, like Bayes, presumes knowledge of a prior probability distribution (pdf) and
12 produces the posterior pdf itself. MRE attempts to produce this pdf based on the
13 information provided by new moments. It will use moments of the prior distribution
14 only if new data on these moments is not available. It is important to note that MRE
15 makes a strong statement that the imposed constraints are exact and complete.
16 In this way, MRE is maximally uncommitted with respect to unknown information.
17 In general, since input data are known only to within a certain accuracy, it is im-
18 portant that any inversion method should allow for errors in the measured data.
19 The MRE approach can accommodate such uncertainty and in new work described
20 here, previous results are modified to include a Gaussian prior. A variety of MRE
21 solutions are reproduced under a number of assumed moments and these include
22 second-order central moments. Various solutions of Jacobs and van der Geest (1991)
23 were repeated and clarified. Menke’s weighted minimum length solution was shown

24 to have a basis in information theory, and the classic least squares estimate is shown
 25 as a solution to MRE under the conditions of more data than unknowns and where
 26 we utilize the observed data and their associated noise. An example inverse problem
 27 involving a gravity survey over a layered and faulted zone is shown. In all cases the
 28 inverse results match quite closely the actual density profile, at least in the upper
 29 portions of the profile. The similar results to Bayes presented in are a reflection of
 30 the fact that the MRE posterior pdf, and its mean are constrained not by $\mathbf{d} = \mathbf{Gm}$
 31 but by its first moment $E(\mathbf{d} = \mathbf{Gm})$, a weakened form of the constraints. If there
 32 is no error in the data then one should expect a complete agreement between Bayes
 33 and MRE and this is what is shown. Similar results are shown when second moment
 34 data is available (for example posterior covariance equal to zero). But dissimilar
 35 results are noted when we attempt to derive a Bayesian like result from MRE. In
 36 the various examples given in this paper, the problems look similar but are, in the
 37 final analysis, not equal. The methods of attack are different and so are the results
 38 even though we have used the linear inverse problem as a common template.

39 1 INTRODUCTION

40 Arguably, since the advent of Taratola's work (eg. Tarantola, 1987) it has become common to
 41 consider both measured data and unknown model parameters as uncertain. In a probabilis-
 42 tic inverse approach, the ultimate goal is the posterior probability density function (pdf),
 43 updated from some previous level of knowledge. Generally, although not exclusively, we are
 44 concerned with the expected values (or some other mode) of the posterior pdf, together with
 45 appropriate confidence limits. A detailed discussion of information-based methods is given in
 46 Ulrych and Sacchi (2006), specifically Bayesian inference, maximum entropy, and minimum
 47 relative entropy (MRE).

48 Entropy maximization (MaxEnt) is a general approach of inferring a probability distribu-
 49 tion from constraints which do not uniquely characterize that distribution. Applications of this
 50 method have met with considerable success in a variety of fields (eg. Kapur 1989; Buck and
 51 Macaulay, 1991). An often studied use for an entropy measure is as a penalty term or "norm".

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52 In this way the entropy of a model can be maximized subject to non-linear constraints imposed
53 the observed data.

54 The related, but more general entropic principle is that of minimum relative entropy
55 (MRE). An axiomatic foundation for MRE has been given by Shore and Johnson (1980)
56 and Johnson and Shore (1983). The MRE principle is perhaps less well known than MaxEnt
57 and was first introduced by Kullback (1959) as a method of statistical inference. The MRE
58 approach to inversion was originally developed by Shore (1981) as an extension of Burg's
59 (1975) method of maximum entropy spectral analysis (see also Ulrych and Bishop, 1975).

60 In the hydrological and geophysical literature our first publications on MRE attracted a
61 certain amount of attention (Woodbury and Ulrych, 1993; 1996; Woodbury et al., 1998a).
62 Our (1998) work was followed up with studies of small 'toy' problems, die experiments and
63 an exhaustive comparison of MRE with Maximum Entropy, Bayesian and SVD solutions
64 (Woodbury and Ulrych, 1998b). Finally, thesis and other papers (Neupauer, 1999; Neupauer et
65 al., 2000; Ulrych and Woodbury, 2003) were produced in which detailed comparisons between
66 Tikhonov regularization (Provencher, 1982ab; TR) and MRE were made.

67 Ulrych et al. (2000) used entropic principles for a classical problem; that is finding moments
68 and distributions from sample data. They showed that probability density functions that are
69 estimated from L-moments are superior estimates to those obtained using sample central
70 moments (C-moments) and the principle of maximum entropy. Woodbury (2004) detailed a
71 general purpose computer program that produces a univariate pdf from a series of constraints
72 and a prior probability. Some guidelines for the selection of the prior were presented.

73 From a Bayesian perspective, Ulrych et al. (2001) summarized some of the concepts central
74 to that approach to inverse problems. Of course, there are many examples in the geophysical
75 literature on Bayes, many of which are referred to in Ulrych and Sacchi (2006) and Scales et.
76 al. (2001). Some of the most often referred to works include Duijndam (1988) and Bretthorst
77 (1988). These references are by no means exhaustive and only serve the purpose of recognizing
78 the importance of Bayes theorem in problems of inference.

79 Empirical Bayesian techniques have also undergone development and application in Geo-
80 physics (e.g. Woodbury, 2007). There has always been controversy surrounding the use of,
81 perhaps, arbitrary priors in Bayes Theorem. For example, the form of the pdf may be known
82 or highly suspected but the actual statistical parameters embedded into the prior pdf, such
83 as the mean, the variance and so on may not be well known and difficult to estimate. Noise
84 levels in the data may also be unknown. The idea behind the empirical Bayes approach is
85 that the prior is based on information contained in the input data (Ulrych et al., 2001).

86 What is perhaps less obvious from the published literature is how these methods are
 87 distinct from classical methods of statistical inference, such as maximum likelihood or least
 88 squares. How do Bayesian methods compare to entropic techniques, and under what circum-
 89 stances should one use entropic rather than Bayesian methods? This current effort is organized
 90 as follows. Minimum relative entropy (MRE), and Bayesian solutions for inverse problems are
 91 detailed and related. The focus of the paper is to illustrate important philosophies on in-
 92 version and the similarities and differences between the various solution methods. While there
 93 exists a considerable body of published works comparing maximum entropy and Bayes, there
 94 are few such works related to Bayes and MRE. Much of the existing work can be attributed
 95 to Kapur and coauthors, but some of the derivations were left incomplete. This paper will
 96 attempt to fill in those gaps, step by step in comparing inverse solutions on the general linear
 97 underdetermined model.

98 2 BAYESIAN APPROACH TO INVERSION

99 In this section we will review the classic Bayesian solution to the linear inverse problem. Much
 100 of this material will likely be well known to the reader but it is appropriate here to repeat this
 101 for the sake of completeness and understanding of the notation used. Here, we rely on linear
 102 transformations and Gaussian assumptions for probabilities. This step makes the following
 103 expectations and analysis tractable.

104 In the discrete linear-inverse case, one needs to start with a forward modeling problem
 105 and this can be written as

$$106 \quad \mathbf{d} = \mathbf{G}\mathbf{m} \tag{1}$$

107 where \mathbf{d} is a $(n \times 1)$ vector of theoretically predicted data and \mathbf{G} $(m \times n)$ is a linear trans-
 108 formation (kernel matrix) from model to data space. In the case of observations though, the
 109 true “data” are corrupted by noise, for example

$$110 \quad \mathbf{d}^* = \mathbf{G}\mathbf{m} + \mathbf{e} \tag{2}$$

111 where \mathbf{e} is an $(n \times 1)$ vector of unknown error terms and \mathbf{m} $(m \times 1)$ vector of model parameters,
 112 “the model”. In the case where there are more data than unknowns we can, of course, use the
 113 classic least-squares approach, but this subject will not be covered in the present work. Here we
 114 focus on more interesting cases where more unknowns are sought than data values observed.
 115 Consequently, this ill-posed deterministic problem is recast in terms of a problem in statistical
 116 inference; one in which can be attacked from the viewpoint of updating probabilities.

117 Bayesian inference supposes that an observer can define a personal prior probability-
 118 density function (pdf) about some random variable \mathbf{m} . This pdf, $p(\mathbf{m})$, can be defined on the
 119 basis of personal experience or judgment and Bayes' rule quantifies how this personal pdf can
 120 be changed on the basis of measurements. Consider a vector of observed random variables,
 121 \mathbf{d}^* . If the conditional pdf of \mathbf{d}^* given the "true" value of \mathbf{m} is given by $\Phi(\mathbf{d}^* | \mathbf{m})$ then
 122 Bayes' rule states that

$$123 \quad \Phi(\mathbf{m} | \mathbf{d}^*) = \frac{\Phi(\mathbf{d}^* | \mathbf{m})p(\mathbf{m})}{\int_K \Phi(\mathbf{d}^* | \mathbf{m})p(\mathbf{m})d\mathbf{m}} \quad (3)$$

124 $\Phi(\mathbf{m} | \mathbf{d}^*)$ is the conditional pdf of \mathbf{m} , given \mathbf{d}^* , and $\Phi(\mathbf{d}^* | \mathbf{m})$ represents the conditional
 125 pdf from forward modeling. If a likelihood function can be defined (i.e., the forward model
 126 exists), and there is a compatibility between observed results and a prior understanding of
 127 model parameters, (i.e., $\Phi(\mathbf{d}^* | \mathbf{m}) > 0$ for some \mathbf{m} where $p(\mathbf{m}) > 0$) then Bayes' rule implies
 128 that the resulting posterior pdf exists and is unique (Tarantola, 1987, p53).

129 For Bayes we must define a likelihood function which is a conditional pdf of the data, given
 130 the "true" model \mathbf{m} . In other words we need a pdf for $(\mathbf{d}^* - \mathbf{d} = \mathbf{e})$. If the physics is correct then
 131 the pdf for \mathbf{e} only reflects measurement error. Classically, it is assumed that $\mathbf{e} = N(\mathbf{0}, \mathbf{C}_d)$; i.e.
 132 normally distributed with mean zero and covariance $\mathbf{C}_d = E[(\mathbf{d} - \mathbf{d}^*)(\mathbf{d} - \mathbf{d}^*)^T]$. Therefore,
 133 we can write the Bayes likelihood as (see Tarantola, 1987, p68):

$$134 \quad \Phi(\mathbf{d}^* | \mathbf{m}) = ((2\pi)^{nd}|\mathbf{C}_d|)^{-\frac{1}{2}} \exp \left[-\frac{1}{2}(\mathbf{d}^* - \mathbf{G}\mathbf{m})^T \mathbf{C}_d^{-1}(\mathbf{d}^* - \mathbf{G}\mathbf{m}) \right] \quad (4)$$

135 where n is the length of vector \mathbf{d}^* . If the prior distribution of the model is also assumed to
 136 be Gaussian then the prior pdf for \mathbf{m} is given the form:

$$137 \quad p(\mathbf{m}) = ((2\pi)^{nm}|\mathbf{C}_p|)^{-\frac{1}{2}} \exp \left[-\frac{1}{2}(\mathbf{m} - \mathbf{s})^T \mathbf{C}_p^{-1}(\mathbf{m} - \mathbf{s}) \right] \quad (5)$$

138 Here, \mathbf{s} is the mean value and \mathbf{C}_p is the covariance matrix which of course describes the
 139 variance and correlation of the parameters. Using (3, 4, and 5) the resulting posterior pdf
 140 from a Bayes analysis of the linear inverse problem with Gaussian priors and likelihood is:

$$\begin{aligned} \Phi(\mathbf{m} | \mathbf{d}^*) &= C_1 \exp \left[-\frac{1}{2}(\mathbf{d}^* - \mathbf{G}\mathbf{m})^T \mathbf{C}_d^{-1}(\mathbf{d}^* - \mathbf{G}\mathbf{m}) \right. \\ &\quad \left. -\frac{1}{2}(\mathbf{m} - \mathbf{s})^T \mathbf{C}_p^{-1}(\mathbf{m} - \mathbf{s}) \right] \end{aligned} \quad (6)$$

141 Note that the mode of a Gaussian distribution is equal to its mean. So to find the mean
 142 and covariance of this pdf, we simply have to find the maximum of an objective function
 143 defined as

$$J(\mathbf{m}) = -\frac{1}{2}(\mathbf{d}^* - \mathbf{G}\mathbf{m})^T \mathbf{C}_d^{-1}(\mathbf{d}^* - \mathbf{G}\mathbf{m})$$

$$-\frac{1}{2}(\mathbf{m} - \mathbf{s})^T \mathbf{C}_p^{-1} (\mathbf{m} - \mathbf{s}) \quad (7)$$

144 Expanding out the above results in

$$\begin{aligned} J(\mathbf{m}) = & -\frac{1}{2}[\mathbf{m}^T(\mathbf{C}_p^{-1} + \mathbf{G}^T \mathbf{C}_d^{-1} \mathbf{G})\mathbf{m} - \mathbf{m}^T(\mathbf{C}_p^{-1} \mathbf{s} + \\ & \mathbf{G}^T \mathbf{C}_d^{-1} \mathbf{d}^*) - (\mathbf{s}^T \mathbf{C}_p^{-1} + \mathbf{d}^{*T} \mathbf{C}_d^{-1} \mathbf{G})\mathbf{m} + \\ & \mathbf{s}^T \mathbf{C}_p^{-1} \mathbf{s} + \mathbf{d}^{*T} \mathbf{C}_d^{-1} \mathbf{d}^*] \end{aligned} \quad (8)$$

145 Letting

$$146 (\mathbf{C}_p^{-1} + \mathbf{G}^T \mathbf{C}_d^{-1} \mathbf{G}) = \mathbf{C}_q^{-1} = \mathbf{A} \quad (9)$$

147 and

$$148 (\mathbf{C}_p^{-1} \mathbf{s} + \mathbf{G}^T \mathbf{C}_d^{-1} \mathbf{d}^*) = \mathbf{b} \quad (10)$$

149 Results in

$$\begin{aligned} J(\mathbf{m}) = & -\frac{1}{2}[\mathbf{m}^T \mathbf{A} \mathbf{m} - \mathbf{m}^T \mathbf{b} - \mathbf{b}^T \mathbf{m} + \mathbf{s}^T \mathbf{C}_p^{-1} \mathbf{s} + \\ & \mathbf{d}^{*T} \mathbf{C}_d^{-1} \mathbf{d}^*] \end{aligned} \quad (11)$$

150 Now taking the derivative of J with respect to \mathbf{m} , recognizing that \mathbf{A} is symmetric and
151 setting the result to zero determines $\langle \mathbf{m} \rangle$, the posterior mean value;

$$152 \mathbf{A} \langle \mathbf{m} \rangle = \mathbf{b} \quad (12)$$

$$153 \langle \mathbf{m} \rangle = \mathbf{A}^{-1} \mathbf{b} \quad (13)$$

154 Expanding,

$$155 \langle \mathbf{m} \rangle = (\mathbf{C}_p^{-1} + \mathbf{G}^T \mathbf{C}_d^{-1} \mathbf{G})^{-1} (\mathbf{C}_p^{-1} \mathbf{s} + \mathbf{G}^T \mathbf{C}_d^{-1} \mathbf{d}^*) \quad (14)$$

$$156 \mathbf{C}_q = (\mathbf{C}_p^{-1} + \mathbf{G}^T \mathbf{C}_d^{-1} \mathbf{G})^{-1} \quad (15)$$

157 which also can be written as:

$$158 \langle \mathbf{m} \rangle = \mathbf{s} + \mathbf{C}_p \mathbf{G}^T (\mathbf{G} \mathbf{C}_p \mathbf{G}^T + \mathbf{C}_d)^{-1} (\mathbf{d}^* - \mathbf{G} \mathbf{s}) \quad (16)$$

$$159 \mathbf{C}_q = \mathbf{C}_p - \mathbf{C}_p \mathbf{G}^T (\mathbf{G} \mathbf{C}_p \mathbf{G}^T + \mathbf{C}_d)^{-1} \mathbf{G} \mathbf{C}_p \quad (17)$$

160 Here, $\langle \mathbf{m} \rangle$ and \mathbf{C}_q are expected value and covariance of the posterior pdf,

$$\begin{aligned} \Phi(\mathbf{m} | \mathbf{d}^*) = & ((2\pi)^{nm} |\mathbf{C}_q|)^{-\frac{1}{2}} \\ & \exp \left[-\frac{1}{2} (\mathbf{m} - \langle \mathbf{m} \rangle)^T \mathbf{C}_q^{-1} (\mathbf{m} - \langle \mathbf{m} \rangle) \right] \end{aligned} \quad (18)$$

161 These results are well known (see Tarantola, 1987). Note in this case $\mathbf{d}^* \neq \mathbf{G} \langle \mathbf{m} \rangle$. In
 162 other words, the average value of the predicted data $\bar{\mathbf{d}}$ from the average value of the model
 163 $\langle \mathbf{m} \rangle$ is not equal to the observed data \mathbf{d}^* . We note that the expected value of the model
 164 (16) could also be written as

$$165 \quad \langle \mathbf{m} \rangle = \mathbf{s} + \mathbf{G}_*^{-1}(\mathbf{d}^* - \mathbf{G}\mathbf{s}) \quad (19)$$

166 where $\mathbf{G}_*^{-1} = \mathbf{C}_p \mathbf{G}^T (\mathbf{G} \mathbf{C}_p \mathbf{G}^T + \mathbf{C}_d)^{-1}$, a generalized inverse matrix of \mathbf{G} .

167 **3 MINIMUM RELATIVE ENTROPY THEORY**

168 Minimum relative entropy (MRE) is an information-theoretic method of problem solving. Its
 169 roots lie in probability theory and in an abstract way deals with information measures in
 170 probability spaces (Kapur and Kessavan, 1992). MRE was classically derived by Shore and
 171 Johnson (1980) and follows directly from four basic axioms of consistent inference which are:
 172 uniqueness, invariance, system independence and subset independence. Stated in other words,
 173 if a problem can be solved in difference ways, each path has to lead to the same answer.

174 Consider a system having a set of possible states. Let \mathbf{x} be a state and $q^\dagger(\mathbf{x})$ its un-
 175 known multivariate probability density function (pdf). Note that $\mathbf{x}^T = (x_1, x_2, x_3, \dots, x_K)$
 176 and the integrals noted below are multiple integrals over each of the x_i . The pdf must satisfy
 177 a normalizing constraint

$$178 \quad \int q^\dagger(\mathbf{x}) d\mathbf{x} = 1 \quad (20)$$

179 Now, let us suppose that there is prior information on $q^\dagger(\mathbf{x})$ in the form of a pdf, $p(\mathbf{x})$, and
 180 new information exists, say expectations of the form of $j = 1, M$ expected value constraints

$$181 \quad \int q^\dagger(\mathbf{x}) f_j(\mathbf{x}) d\mathbf{x} = \bar{f}_j \quad (21)$$

182 or bounds on these values. It is important to note here that MRE makes a strong statement
 183 that the constraints are exact and complete. That is, the M values are the only only ones
 184 operating and these are known exactly. The task at hand then is to choose a distribution $q(\mathbf{x})$
 185 out of the infinite possibilities that is in some way the best estimate of q^\dagger , consistent with this
 186 new information. The solution is to minimize $H(q, p)$, the entropy of $q(\mathbf{x})$ relative to $p(\mathbf{x})$,
 187 where

$$188 \quad H(q, p) = \int q(\mathbf{x}) \ln \left[\frac{q(\mathbf{x})}{p(\mathbf{x})} \right] d\mathbf{x} \quad (22)$$

189 subject to the constraints of the form of equations (20) and (21). The posterior estimate $q(\mathbf{x})$
 190 has the form (Shore and Johnson, 1980; Woodbury and Ulrych, 1993)

$$q(\mathbf{x}) = p(\mathbf{x}) \exp \left[-1 - \mu - \sum_{j=1}^M \lambda_j f_j(\mathbf{x}) \right] \quad (23)$$

where μ and the λ_j are Lagrange multipliers determined from equations (20) and (21). If the final goal is the posterior pdf, (23) is the form of the solution and one has to determine the Lagrange multipliers (Johnson, 1983). For an inverse problem, we are likely concerned with the expected values of the posterior pdf, $q(\mathbf{x})$ together with the appropriate confidence limits.

Now for the geophysical inverse problem, we have to minimize

$$H = \int q(\mathbf{m}) \ln \frac{q(\mathbf{m})}{p(\mathbf{m})} d\mathbf{m} \quad (24)$$

subject to any number of moments but for the sake of argument here, the first two central moments:

$$\int q(\mathbf{m}) \mathbf{m} d\mathbf{m} = \langle \mathbf{m} \rangle \quad (25)$$

$$\int q(\mathbf{m}) [(\langle \mathbf{m} \rangle - \mathbf{m})(\langle \mathbf{m} \rangle - \mathbf{m})^T] d\mathbf{m} = \mathbf{C}_q \quad (26)$$

along with the normalizing constraint

$$\int q(\mathbf{m}) d\mathbf{m} = 1 \quad (27)$$

The MRE solution to this problem is (Kapur, 1989)

$$q(\mathbf{m}) = ((2\pi)^{nm} |\mathbf{C}_q|)^{-\frac{1}{2}} \exp \left[-\frac{1}{2} (\mathbf{m} - \langle \mathbf{m} \rangle)^T \mathbf{C}_q^{-1} (\mathbf{m} - \langle \mathbf{m} \rangle) \right] \quad (28)$$

The reader will notice that in order for us to obtain this solution we would have to first specify moments of the posterior pdf, namely \mathbf{C}_q and $\langle \mathbf{m} \rangle$ which we do not normally have.

Instead, given the theoretical relationship $\mathbf{d} = \mathbf{G}\mathbf{m}$, we can rearrange the above constraints into a form that is more convenient for the data we actually observe and then solve the MRE minimization. Rearranging (25) yields

$$\int q(\mathbf{m}) (\bar{\mathbf{d}} - \mathbf{G}\mathbf{m}) d\mathbf{m} = 0 \quad (29)$$

where $\mathbf{G} \langle \mathbf{m} \rangle = \bar{\mathbf{d}}$ and

$$\int q(\mathbf{m}) [(\bar{\mathbf{d}} - \mathbf{G}\mathbf{m})(\bar{\mathbf{d}} - \mathbf{G}\mathbf{m})^T] d\mathbf{m} = \mathbf{R} = \mathbf{G}\mathbf{C}_q\mathbf{G}^T \quad (30)$$

In the above, \mathbf{R} is the covariance of \mathbf{d} , $E[(\mathbf{d} - \bar{\mathbf{d}})(\mathbf{d} - \bar{\mathbf{d}})^T]$. If $p(\mathbf{m})$ is multivariate and Gaussian;

$$p(\mathbf{m}) = ((2\pi)^m |\mathbf{C}_p|)^{-\frac{1}{2}}$$

$$\exp \left[-\frac{1}{2}(\mathbf{m} - \mathbf{s})^T \mathbf{C}_p^{-1}(\mathbf{m} - \mathbf{s}) \right] \quad (31)$$

215 then $q(\mathbf{m})$ is (Kapur et al., 1994):

$$\begin{aligned} q(\mathbf{m}) &= C_2 \times \exp \left[-\frac{1}{2}((\mathbf{m} - \mathbf{s})^T \mathbf{C}_p^{-1}(\mathbf{m} - \mathbf{s})) \right] \\ &\quad \times \exp \left[-\lambda_0 - \boldsymbol{\lambda}^T(\bar{\mathbf{d}} - \mathbf{G}\mathbf{m}) - (\bar{\mathbf{d}} - \mathbf{G}\mathbf{m})^T \mathbf{D}^{-1}(\bar{\mathbf{d}} - \mathbf{G}\mathbf{m}) \right] \end{aligned} \quad (32)$$

216 where C_2 is a normalizing constant, and λ_0 , $\boldsymbol{\lambda}$ and the matrix \mathbf{D}^{-1} are Lagrange multipliers
217 that have to be determined from the constraints (27, 29, 30). In the subsequent sections it
218 will be shown how the multipliers can be determined for specific cases.

219 **3.1 Updating prior with only first moment information**

220 For the MRE solution, we want to minimize

$$221 \int q(\mathbf{m}) \ln \frac{q(\mathbf{m})}{p(\mathbf{m})} d\mathbf{m} \quad (33)$$

222 subject to:

$$223 \int q(\mathbf{m}) d\mathbf{m} = 1 \quad (34)$$

224 along with

$$225 \int q(\mathbf{m})(\bar{\mathbf{d}} - \mathbf{G}\mathbf{m}) d\mathbf{m} = 0 \quad (35)$$

226 That is, we want to update the prior, based only on a new mean value constraint. As before we
227 assume a Gaussian prior, and replacing the expected value of the data $\bar{\mathbf{d}}$ for the observations
228 \mathbf{d}^* that we actually obtain yields

$$\begin{aligned} q(\mathbf{m}) &= C_3 \times \exp \left[-\frac{1}{2}((\mathbf{m} - \mathbf{s})^T \mathbf{C}_p^{-1}(\mathbf{m} - \mathbf{s})) \right] \\ &\quad \times \exp \left[-\lambda_0 - \boldsymbol{\lambda}^T(\mathbf{d}^* - \mathbf{G}\mathbf{m}) \right] \end{aligned} \quad (36)$$

229 Defining $\mathbf{A} = \mathbf{C}_p^{-1}$ and $\mathbf{b} = \mathbf{C}_p^{-1}\mathbf{s}$ yields the objective function

$$\begin{aligned} J(\mathbf{m}) &= -\frac{1}{2}[\mathbf{m}^T \mathbf{A}\mathbf{m} - \mathbf{m}^T \mathbf{b} - \mathbf{b}^T \mathbf{m} + \mathbf{s}^T \mathbf{C}_p^{-1} \mathbf{s} + \\ &\quad -2\boldsymbol{\lambda}^T (\mathbf{G}\mathbf{m} - \mathbf{d}^*)] \end{aligned} \quad (37)$$

230 Taking the derivative of J with respect to \mathbf{m} and setting the result to zero yields

$$231 \mathbf{A} \langle \mathbf{m} \rangle = \mathbf{b} - \mathbf{G}^T \boldsymbol{\lambda} \quad (38)$$

232 The next step is to find the λ values. We know from the second MRE constraint

$$233 \int q(\mathbf{m})(\mathbf{d}^* - \mathbf{G}\mathbf{m})d\mathbf{m} = 0 \quad (39)$$

234 or $\mathbf{d}^* = \mathbf{G} \langle \mathbf{m} \rangle$. Substituting the above expression for the mean value of the model results
235 in

$$236 \mathbf{G}\mathbf{A}^{-1}(\mathbf{b} + \mathbf{G}^T\lambda) = \mathbf{d}^* \quad (40)$$

237 and solving for the λ 's

$$238 \lambda = (\mathbf{G}\mathbf{A}^{-1}\mathbf{G}^T)^{-1}[\mathbf{d}^* - \mathbf{G}\mathbf{A}^{-1}\mathbf{b}] \quad (41)$$

239 Finally substituting this value back into our expression for the mean value

$$240 \langle \mathbf{m} \rangle = \mathbf{A}^{-1} [\mathbf{b} + \mathbf{G}^T\lambda] \quad (42)$$

$$241 \langle \mathbf{m} \rangle = \mathbf{A}^{-1}(\mathbf{b} + \mathbf{G}^T(\mathbf{G}\mathbf{A}^{-1}\mathbf{G}^T)^{-1}[\mathbf{d}^* - \mathbf{G}\mathbf{A}^{-1}\mathbf{b}]) \quad (43)$$

$$242 \langle \mathbf{m} \rangle = \mathbf{A}^{-1}\mathbf{b} + \mathbf{A}^{-1}\mathbf{G}^T(\mathbf{G}\mathbf{A}^{-1}\mathbf{G}^T)^{-1}[\mathbf{d}^* - \mathbf{G}\mathbf{A}^{-1}\mathbf{b}] \quad (44)$$

243 or, with the definitions of $\mathbf{A} = \mathbf{C}_p^{-1}$ and $\mathbf{b} = \mathbf{C}_p^{-1}\mathbf{s}$

$$244 \langle \mathbf{m} \rangle = \mathbf{s} + \mathbf{C}_p\mathbf{G}^T(\mathbf{G}\mathbf{C}_p\mathbf{G}^T)^{-1}[\mathbf{d}^* - \mathbf{G}\mathbf{s}] \quad (45)$$

245 and it is noted that the posterior covariance is not updated;

$$246 \mathbf{C}_q = \mathbf{C}_p \quad (46)$$

247 This is identical to the results of Jacobs and van der Geest (1991). Note that $\mathbf{C}_p\mathbf{G}^T(\mathbf{G}\mathbf{C}_p\mathbf{G}^T)^{-1}$
248 could be again be referred to as \mathbf{G}_*^{-1} , a generalized inverse of \mathbf{G} . Equation (45) is of the same
249 form as a weighted minimum length solution (see Menke, 1989, p. 54). Therefore, the WML
250 solution has a basis in information theory.

251 3.2 Updating prior with first and second moment information

252 For the MRE solution, we have to minimize (33) subject to (34) along with the mean value
253 constraint (35) and

$$254 \int q(\mathbf{m})[(\bar{\mathbf{d}} - \mathbf{G}\mathbf{m})(\bar{\mathbf{d}} - \mathbf{G}\mathbf{m})^T]d\mathbf{m} = \mathbf{R} \quad (47)$$

255 $q(\mathbf{m})$ is of the form (32) and again, the Lagrange multipliers determined from the con-
256 straints. We can then proceed in the same way as before;

257 Letting

$$258 (\mathbf{C}_p^{-1} + \mathbf{G}^T\mathbf{D}^{-1}\mathbf{G}) = \mathbf{C}_q^{-1} = \mathbf{A} \quad (48)$$

259 and

$$260 \quad (\mathbf{C}_p^{-1} \mathbf{s} + \mathbf{G}^T \mathbf{D}^{-1} \bar{\mathbf{d}}) = \mathbf{b} \quad (49)$$

261 We now have

$$J(\mathbf{m}) = -\frac{1}{2} [\mathbf{m}^T \mathbf{A} \mathbf{m} - \mathbf{m}^T \mathbf{b} - \mathbf{b}^T \mathbf{m} + \mathbf{s}^T \mathbf{C}_p^{-1} \mathbf{s} + \bar{\mathbf{d}}^T \mathbf{D}^{-1} \bar{\mathbf{d}} - 2\lambda^T (\mathbf{G} \mathbf{m} - \bar{\mathbf{d}})] \quad (50)$$

262 Now taking the derivative of J , and setting the result to zero determines $\langle \mathbf{m} \rangle$, as

$$263 \quad \mathbf{A} \langle \mathbf{m} \rangle - \mathbf{b} - \lambda^T \mathbf{G} = 0 \quad (51)$$

$$264 \quad \mathbf{A} \langle \mathbf{m} \rangle = \mathbf{b} + \lambda^T \mathbf{G} \quad (52)$$

$$265 \quad \langle \mathbf{m} \rangle = \mathbf{A}^{-1} [\mathbf{b} + \lambda^T \mathbf{G}] \quad (53)$$

266 Expanding these terms out yields

$$267 \quad \langle \mathbf{m} \rangle = (\mathbf{C}_p^{-1} + \mathbf{G}^T \mathbf{D}^{-1} \mathbf{G})^{-1} (\mathbf{C}_p^{-1} \mathbf{s} + \mathbf{G}^T \mathbf{D}^{-1} \bar{\mathbf{d}} - \mathbf{G}^T \lambda) \quad (54)$$

$$268 \quad \mathbf{C}_q = (\mathbf{C}_p^{-1} + \mathbf{G}^T \mathbf{D}^{-1} \mathbf{G})^{-1} \quad (55)$$

269 Rearranging,

$$270 \quad \langle \mathbf{m} \rangle = \mathbf{C}_q (\mathbf{C}_p^{-1} \mathbf{s} + \mathbf{G}^T \mathbf{D}^{-1} \bar{\mathbf{d}} - \mathbf{G}^T \lambda) \quad (56)$$

271 The covariance can be written as

$$272 \quad \mathbf{C}_q = \mathbf{C}_p - \mathbf{C}_p \mathbf{G}^T [\mathbf{G} \mathbf{C}_p \mathbf{G}^T + \mathbf{D}]^{-1} \mathbf{G} \mathbf{C}_p \quad (57)$$

273 Let us first examine a limiting case. Suppose we impose a second moment constraint
274 condition that $\mathbf{C}_q = 0$. This means that $\mathbf{R} = 0$, and then \mathbf{D} must also equal zero in the above
275 and

$$276 \quad \mathbf{C}_q = \mathbf{C}_p - \mathbf{C}_p \mathbf{G}^T (\mathbf{G} \mathbf{C}_p \mathbf{G}^T)^{-1} \mathbf{G} \mathbf{C}_p \rightarrow 0 \quad (58)$$

277 Using the results from (44), we also arrive with

$$278 \quad \langle \mathbf{m} \rangle = \mathbf{s} + \mathbf{C}_p \mathbf{G}^T (\mathbf{G} \mathbf{C}_p \mathbf{G}^T)^{-1} [\mathbf{d}^* - \mathbf{G} \mathbf{s}] \quad (59)$$

279 which is identical to the results of Jacob and van Der Geest (1991) for $\mathbf{R} = 0$, assuming
280 $\bar{\mathbf{d}} = \mathbf{d}^*$. In (58) $\mathbf{C}_p \mathbf{G}^T (\mathbf{G} \mathbf{C}_p \mathbf{G}^T)^{-1} = \mathbf{G}_*^{-1}$. The important lesson to be learned at this stage
281 is that the second moment information in the posterior is not updated unless specifically
282 required.

283 3.2.1 General Case: Determining Mean and Covariance

284 Noting the second MRE moment constraint is

$$E[(\mathbf{m} - \langle \mathbf{m} \rangle)(\mathbf{m} - \langle \mathbf{m} \rangle)^T] = \mathbf{C}_q \quad (60)$$

$$\int q(\mathbf{m})(\mathbf{G}\mathbf{m} - \bar{\mathbf{d}})(\mathbf{G}\mathbf{m} - \bar{\mathbf{d}})^T d\mathbf{m} = \mathbf{G}\mathbf{C}_q\mathbf{G}^T = \mathbf{R} \quad (61)$$

The next step is to find the λ values. We know from the first MRE constraint.

$$\int q(\mathbf{m})(\bar{\mathbf{d}} - \mathbf{G}\mathbf{m})d\mathbf{m} = 0 \quad (62)$$

or $\bar{\mathbf{d}} = \mathbf{G} \langle \mathbf{m} \rangle$. Substituting the above expression for the mean value of the model (58) results in

$$\langle \mathbf{m} \rangle = \mathbf{C}_q(\mathbf{C}_p^{-1}\mathbf{s} + \mathbf{G}^T\mathbf{D}^{-1} - \bar{\mathbf{d}}^T\lambda) \quad (63)$$

$$\langle \mathbf{m} \rangle = (\mathbf{C}_q\mathbf{C}_p^{-1}\mathbf{s} + \mathbf{C}_q\mathbf{G}^T\mathbf{D}^{-1} - \bar{\mathbf{d}}_q\mathbf{G}^T\lambda) \quad (64)$$

Now we know from the mean value constraint that $\mathbf{G} \langle \mathbf{m} \rangle = \bar{\mathbf{d}}$ and multiplying the mean model by \mathbf{G} yields

$$\mathbf{G} \langle \mathbf{m} \rangle = \bar{\mathbf{d}} = \mathbf{G}\mathbf{C}_q\mathbf{C}_p^{-1}\mathbf{s} + \mathbf{G}\mathbf{C}_q\mathbf{G}^T\mathbf{D}^{-1}\bar{\mathbf{d}} - \mathbf{G}\mathbf{C}_q\mathbf{G}^T\lambda \quad (65)$$

Solving for λ and substituting back into the expression for the mean value results in

$$\begin{aligned} \langle \mathbf{m} \rangle &= \mathbf{C}_q\mathbf{C}_p^{-1}\mathbf{s} + \mathbf{C}_q\mathbf{G}^T\mathbf{D}^{-1} + \mathbf{C}_q\mathbf{G}^T\mathbf{R}^{-1}\bar{\mathbf{d}} \\ &\quad - \mathbf{C}_q\mathbf{G}^T[\mathbf{G}\mathbf{C}_q\mathbf{G}^T]^{-1}\mathbf{G}\mathbf{C}_q\mathbf{C}_p^{-1}\mathbf{s} \\ &\quad - \mathbf{C}_q\mathbf{G}^T\mathbf{D}^{-1}\bar{\mathbf{d}} \end{aligned} \quad (66)$$

This equation may simplify to

$$\langle \mathbf{m} \rangle = \mathbf{C}_q\mathbf{C}_p^{-1}\mathbf{s} + \mathbf{C}_q\mathbf{G}^T\mathbf{R}^{-1}\bar{\mathbf{d}} - \mathbf{C}_q\mathbf{C}_p^{-1}\mathbf{s} \quad (67)$$

and finally to

$$\langle \mathbf{m} \rangle = \mathbf{C}_q\mathbf{G}^T\mathbf{R}^{-1}\bar{\mathbf{d}} \quad (68)$$

One can see here that knowledge of \mathbf{C}_q is required to obtain a solution and for many problems this would not be known. If we choose to approximate it as $\mathbf{R} = \mathbf{C}_d = \mathbf{G}\mathbf{C}_q\mathbf{G}^T$ and $\bar{\mathbf{d}} = \mathbf{d}^*$ then

$$\langle \mathbf{m} \rangle = (\mathbf{G}^T\mathbf{C}_d^{-1}\mathbf{G})^{-1}\mathbf{G}^T\mathbf{C}_d^{-1}\mathbf{d}^* \quad (69)$$

The reader should note that the above result is only valid in a "weak" sense, that is the matrix $\mathbf{C}_q = (\mathbf{G}^T\mathbf{C}_d^{-1}\mathbf{G})^{-1}$ in the above case (69) is likely singular for the underdetermined case and in this situation one would have to rely on the generalized inverse. We also note that replacing \mathbf{R} for \mathbf{C}_d involves an approximation that Bayes does not have. Equation (69) can be recognized as the classic least squares solution if there is more data than unknowns. Note

310 also the prior covariance is ignored. This should perhaps come as no surprise to us that MRE
 311 ignores moments of the prior information if not strictly required.

312 We can write the equation for the mean as

$$313 \quad \langle \mathbf{m} \rangle = \mathbf{G}_*^{-1} \mathbf{d}^* \quad (70)$$

$$314 \quad \text{where } \mathbf{G}_*^{-1} = \mathbf{C}_q \mathbf{G}^T \mathbf{C}_d^{-1}.$$

3.2.2 *Constraining Covariance Only*

316 If we choose to enforce the second moment constraint (63) and not the mean constraint (35),
 317 then we do not require that the data are fitted exactly. By examining equations (58, 59) the
 318 constraints on $\boldsymbol{\lambda}$ are not required. This reduces (58) and (59) to

$$319 \quad \langle \mathbf{m} \rangle = \mathbf{C}_q (\mathbf{C}_p^{-1} \mathbf{s} + \mathbf{G}^T \mathbf{D}^{-1} \bar{\mathbf{d}}) \quad (71)$$

320 and the posterior covariance

$$321 \quad \mathbf{C}_q = \mathbf{C}_p - \mathbf{C}_p \mathbf{G}^T [\mathbf{G} \mathbf{C}_p \mathbf{G}^T + \mathbf{D}]^{-1} \mathbf{G} \mathbf{C}_p \quad (72)$$

322 MRE requires that we know \mathbf{C}_q and then given this we would have to find that matrix \mathbf{D} to
 323 satisfy the constraint. If it is *assumed* that the posterior covariance is of the same form as
 324 Bayes

$$325 \quad \mathbf{C}_q = \mathbf{C}_p - \mathbf{C}_p \mathbf{G}^T (\mathbf{G} \mathbf{C}_p \mathbf{G}^T + \mathbf{C}_d^{-1})^{-1} \mathbf{G} \mathbf{C}_p \quad (73)$$

326 Then clearly $\mathbf{D} = \mathbf{C}_d^{-1}$ in (73) and (74). Using the results from (58), we also arrive with

$$327 \quad \langle \mathbf{m} \rangle = \mathbf{s} + \mathbf{C}_p \mathbf{G}^T (\mathbf{G} \mathbf{C}_p \mathbf{G}^T + \mathbf{C}_d^{-1})^{-1} [\mathbf{d}^* - \mathbf{G} \mathbf{s}] \quad (74)$$

328 assuming $\bar{\mathbf{d}} = \mathbf{d}^*$.

3.3 **Updating prior with first moment information and uncertain constraints**

330 In general, since the input data are known only to within a certain accuracy, it is important
 331 that we do not satisfy the first moment constraint in a strong way, instead $\mathbf{d}^* \neq \mathbf{G} \langle \mathbf{m} \rangle = \bar{\mathbf{d}}$. The inversion method should allow for errors in the measured data. The MRE
 332 approach can accommodate such uncertainty and this subject has been discussed by Johnson
 333 and Shore (1984), Ulrych et al. (1990) and Woodbury and Ulrych (1998b).

335 The problem is posed in the following way. Minimize $H(q, p)$, the entropy of $q(\mathbf{m})$ relative
 336 to $p(\mathbf{m})$, (34), subject to (35) and a more general form of equation (21);

$$337 \quad [(\mathbf{G} \langle \mathbf{m} \rangle - \bar{\mathbf{d}})^T (\mathbf{G} \langle \mathbf{m} \rangle - \bar{\mathbf{d}})] \leq \epsilon^2 \quad (75)$$

338 or

$$339 \left(\int q(\mathbf{m}) \mathbf{G} \mathbf{m} d\mathbf{m} - \mathbf{d}^* \right)^T \left(\int q(\mathbf{m}) \mathbf{G} \mathbf{m} d\mathbf{m} - \mathbf{d}^* \right) \leq \epsilon^2 \quad (76)$$

340 where ϵ is a known error term and replacing $\bar{\mathbf{d}}$ for \mathbf{d}^* . Note that ϵ is required for the analysis to
 341 proceed. If this value is not known, Ulrych and Woodbury (2003) show that it can be estimated
 342 in a variety of ways, including estimated from the data themselves, a rigorous approach using
 343 the real cepstrum and the AIC criterion. It can be shown (Johnson and Shore, 1984) that
 344 the MRE solution has the same form as equations (10) and (11) but with data constraints
 345 modified to

$$346 \mathbf{G} \langle \mathbf{m} \rangle = \mathbf{d}^* - \epsilon \frac{\boldsymbol{\beta}}{\|\boldsymbol{\beta}\|} = \mathbf{d}^\dagger \quad (77)$$

347 for the case where the data errors are identically distributed and independent. The vector
 348 $\boldsymbol{\beta} = (\beta_1, \beta_2, \dots, \beta_n)$ is a vector of Lagrange multipliers. The important fact from the point
 349 of view of MRE is that the form of the solution with uncertain data is not changed; only a
 350 simple modification of the solution algorithm is required.

351 Past efforts in MRE along these lines take the point of view of the prior either being
 352 uniform or with some character such as a truncated exponential (Woodbury and Ulrych,
 353 1996). In new work described here, previous results are modified to include a Gaussian prior.
 354 In this light we can use the following results. First define our posterior pdf in the form

$$355 q(\mathbf{m}) = p(\mathbf{m}) \exp \left[-\alpha - \boldsymbol{\beta}^T \mathbf{G} \mathbf{m} \right] \quad (78)$$

356 where

$$357 \boldsymbol{\beta} = 2\lambda \mathbf{G} \langle \mathbf{m} \rangle \quad (79)$$

358 If $p(\mathbf{m})$ is Gaussian as before then

$$q(\mathbf{m}) = C_3 \times \exp \left[-\frac{1}{2} ((\mathbf{m} - \mathbf{s})^T \mathbf{C}_p^{-1} (\mathbf{m} - \mathbf{s})) \right] \quad (80)$$

$$\times \exp \left[-\lambda_0 - \boldsymbol{\beta}^T (\mathbf{d}^\dagger - \mathbf{G} \mathbf{m}) \right] \quad (81)$$

359 As shown earlier, the mean value for $q(\mathbf{m})$ is

$$360 \langle \mathbf{m} \rangle = \mathbf{s} + \mathbf{C}_p \mathbf{G}^T (\mathbf{G} \mathbf{C}_p \mathbf{G}^T)^{-1} [\mathbf{d}^\dagger - \mathbf{G} \mathbf{s}] \quad (82)$$

361 and we note $\mathbf{C}_q = \mathbf{C}_p$. Note also these results could be expressed in a general way,

$$362 \langle \mathbf{m} \rangle = \mathbf{s} + \mathbf{G}_*^{-1} [\mathbf{d}^\dagger - \mathbf{G} \mathbf{s}] \quad (83)$$

$$363 \boldsymbol{\beta} = (\mathbf{G} \mathbf{C}_p \mathbf{G}^T)^{-1} [\mathbf{d}^\dagger - \mathbf{G} \mathbf{s}] \quad (84)$$

364 The mean value written in terms of $\boldsymbol{\beta}$

$$\langle \mathbf{m} \rangle = \mathbf{C}_p \left[\mathbf{C}_p^{-1} \mathbf{s} + \mathbf{G}^T \boldsymbol{\beta} \right] \quad (85)$$

The above equation for $\langle \mathbf{m} \rangle$ includes a non-linear dependency on the Lagrange multipliers β_j in the term for \mathbf{d}^\dagger and an iterative sequence must be used to establish the final result. Note that this result can be shown to be identical in form to that of the WML solution of (45).

4 COMPARISONS OF BAYESIAN AND MRE APPROACHES

In this section we will summarize the approaches detailed in this paper and offer a comparison between them (see Tables I and II). Seth and Kapur (1990), Kapur and Kessavan (1992), Kapur et al. (1994), Macaulay and Buck (1989) and Jacobs and van der Geest (1991) enunciated some of these differences and the essence of their comparisons is repeated and expanded upon here. The essential difference between the two is given below:

(i) Bayes: This approach is different than the classical methods of first obtaining a sample, and then using sampling distribution theory to obtain estimates of the parameters. Bayes assumes some prior probability (density) which the model follows. The method then proceeds to use a sample of observations, say x_1, x_2, \dots to update the prior probability to a new (revised) posterior probability. This new probability, or probability density, incorporates the information of the sample. More data can be then taken and the updated pdf can be renewed again in a continuous way. These features of Bayes, the assumptions of prior probabilities and continuous updating in the light of new observations, are of key importance. The prior pdf can play an important role in the inversion and has been the center of many disputes that have come about as the result of the use Bayes theorem.

(ii) MRE: This approach differs from both Bayes and classical statistical methods in that knowledge of moments are used as "data" rather than sample values, x_1, x_2, \dots . MRE, like Bayes, presumes knowledge of a prior probability distribution and produces the posterior pdf itself. For example, suppose we have a univariate prior pdf that is Gaussian and we know the mean μ and variance σ^2 . Suppose we have new information in the form of a mean μ_1 . Then the posterior pdf is also Gaussian with mean changed to μ_1 but with the variance unchanged. MRE attempts to produce a pdf based on the information provided by new moments. It will use moments of the prior distribution only if new data on these moments is not available (see Woodbury, 2004). It is important to note here that MRE makes a strong statement that the constraints are exact and complete. That is, the M values are the only ones operating

396 and these are known exactly. In this way, MRE is maximally uncommitted with respect to
 397 unknown information.

398 Some examples of Bayesian inference along with MRE solutions are given below (Tables I
 399 and II). These would appear to be at first glance identical problems and are used to illustrate
 400 important differences and similarities. For the case of a linear inverse problem and the Bayes
 401 solution, imagine that the observations are error free; that is $\mathbf{C}_d \rightarrow \mathbf{0}$ (Table I, row 4). Using
 402 the relationships

$$403 \quad \langle \mathbf{m} \rangle = \mathbf{s} + \mathbf{C}_p \mathbf{G}^T (\mathbf{G} \mathbf{C}_p \mathbf{G}^T + \mathbf{C}_d)^{-1} (\mathbf{d}^* - \mathbf{G} \mathbf{s}) \quad (86)$$

404 with $\mathbf{C}_d = \mathbf{0}$ results in

$$405 \quad \langle \mathbf{m} \rangle = \mathbf{s} + \mathbf{C}_p \mathbf{G}^T (\mathbf{G} \mathbf{C}_p \mathbf{G}^T)^{-1} (\mathbf{d}^* - \mathbf{G} \mathbf{s}) \quad (87)$$

406 This mean value is identical to the MRE approach of updating based on the first moment
 407 only (Table I, row 5). However, the Bayes posterior covariance with $\mathbf{C}_d = \mathbf{0}$ is;

$$408 \quad \mathbf{C}_q = \mathbf{C}_p - \mathbf{C}_p \mathbf{G}^T (\mathbf{G} \mathbf{C}_p \mathbf{G}^T)^{-1} \mathbf{G} \mathbf{C}_p = \mathbf{C}_p - \mathbf{G}_*^{-1} \mathbf{G} \mathbf{C}_p \rightarrow \mathbf{0} \quad (88)$$

409 which is identical to the MRE result (Table I, row 3) but an additional constraint is
 410 required. Let us take the opposite case for the Bayes solution and imagine that the observations
 411 are totally inaccurate; that is $\mathbf{C}_d \rightarrow \infty$ or $\mathbf{C}_d^{-1} \rightarrow \mathbf{0}$. Recall that for Bayes

$$412 \quad \langle \mathbf{m} \rangle = (\mathbf{C}_p^{-1} + \mathbf{G}^T \mathbf{C}_d^{-1} \mathbf{G})^{-1} (\mathbf{C}_p^{-1} \mathbf{s} + \mathbf{G}^T \mathbf{C}_d^{-1} \mathbf{d}^*) \quad (89)$$

413 If $\mathbf{C}_d^{-1} \rightarrow \mathbf{0}$ then $\mathbf{m} = \mathbf{s}$ and the covariance $\mathbf{C}_q = \mathbf{C}_p$. (See Table I, row 6). In this limiting
 414 case if we assume that the data are totally imprecise then the Bayes posterior mean is equal
 415 to the prior mean and the prior covariance is equal to the posterior covariance. These results
 416 are not the same for the MRE approach in the mean (Table I, row 5) but identical in the
 417 covariance. Note though, that not specifying a covariance is not the same as stipulating that
 418 the covariance is infinite.

419 Table I, row 7 shows how we can alter the MRE approach in such a way that it includes
 420 an estimate of the error in the observed data. In this way the results for the mean value are
 421 similar in form to pure Bayes (Table I, row 1), however no change in the covariance is possible
 422 with this approach.

423 Figure 1 shows the result of an example inverse problem involving a gravity survey over
 424 a layered and faulted zone. The example is similar in many respects to that of Mosegaard
 425 and Tarantola (1995). The problem considered is a vertical fault extending from surface to
 426 a maximum depth of 100 km. The example has 26 layers with a thickness derived from an
 427 exponential function with a mean value of 4 km in thickness. The actual layer densities are

428 generated out of an uncorrelated normal distribution with mean value of 3,000 kg/m³ and
 429 a standard deviation of 500 kg/m³. A total 19 values of the first horizontal derivative of
 430 gravity are generated at evenly spaced points (2,000 m) across the fault using equation (1) of
 431 Mohan et al. (1986). These values are corrupted with Gaussian noise of $0.25 \times 10^{-9} \text{ s}^{-2}$. This
 432 information then, forms the kernel matrix and data set in a discrete form.

433 Figure 1 shows in red the actual, or true density distribution with depth. The black line
 434 shows the Bayesian results with a prior Gaussian distribution of standard deviation of 500
 435 kg/m³ and correlation length of 5,000 m. The blue line shows the Bayesian results with the
 436 same prior as above and correlation length of 2,000 m. In green is the equivalent MRE solution
 437 with the same prior as the blue line case above. Here, the inverse matrix in (26) was found
 438 to be singular and an SVD was used to obtain the result. In all cases the inverse results
 439 match quite closely the actual density profile, at least in the upper portions of the profile.
 440 Overall, the general trend in the results is clear and reflects a typical smearing consistent with
 441 expected value determinations. In each case the predicted model generates data that very
 442 closely matches the observations.

443 5 CONCLUSIONS

444 Much of the existing work comparing maximum entropy, MRE and Bayes, can be attributed
 445 to Kapur and coauthors, but some of the derivations in various works were left incomplete.
 446 This paper attempts to fill in those gaps, step by step, in comparing inverse solutions. Specif-
 447 ically an undetermined-discrete linear inverse problem was chosen as a common template for
 448 comparisons. It is assumed here that the “noise” in a set of observables is Gaussian and the
 449 prior probability pdf is also assumed multivariate Gaussian.

450 In the various examples given in this paper, the problems look similar but are, in the final
 451 analysis, not quite equal. The similar results presented in Figure 1 and Table I are a reflection
 452 of the fact that the MRE posterior pdf, and its mean is constrained not by $\mathbf{d} = \mathbf{G}\mathbf{m}$ but by
 453 its first moment $E(\mathbf{d} = \mathbf{G}\mathbf{m})$ a weakened form of the constraints (Jacobs and van der Geest,
 454 1991) . If there is no error in the data then one should expect a complete agreement between
 455 Bayes and MRE and this is what is shown. Similar results are shown when second moment
 456 data is available (for example posterior covariance equal to zero). But dissimilar results are
 457 noted when we attempt to derive a Bayesian like result from MRE (see section 3.2.2). The
 458 MRE pdf is still multivariate Gaussian (see 32) but different than that of Bayes (18) because
 459 that distribution does not satisfy the same constraints. We can derive the Bayes solution from
 460 MRE principles in the undetermined case if we do not insist that the predicted data $\bar{\mathbf{d}}$ are

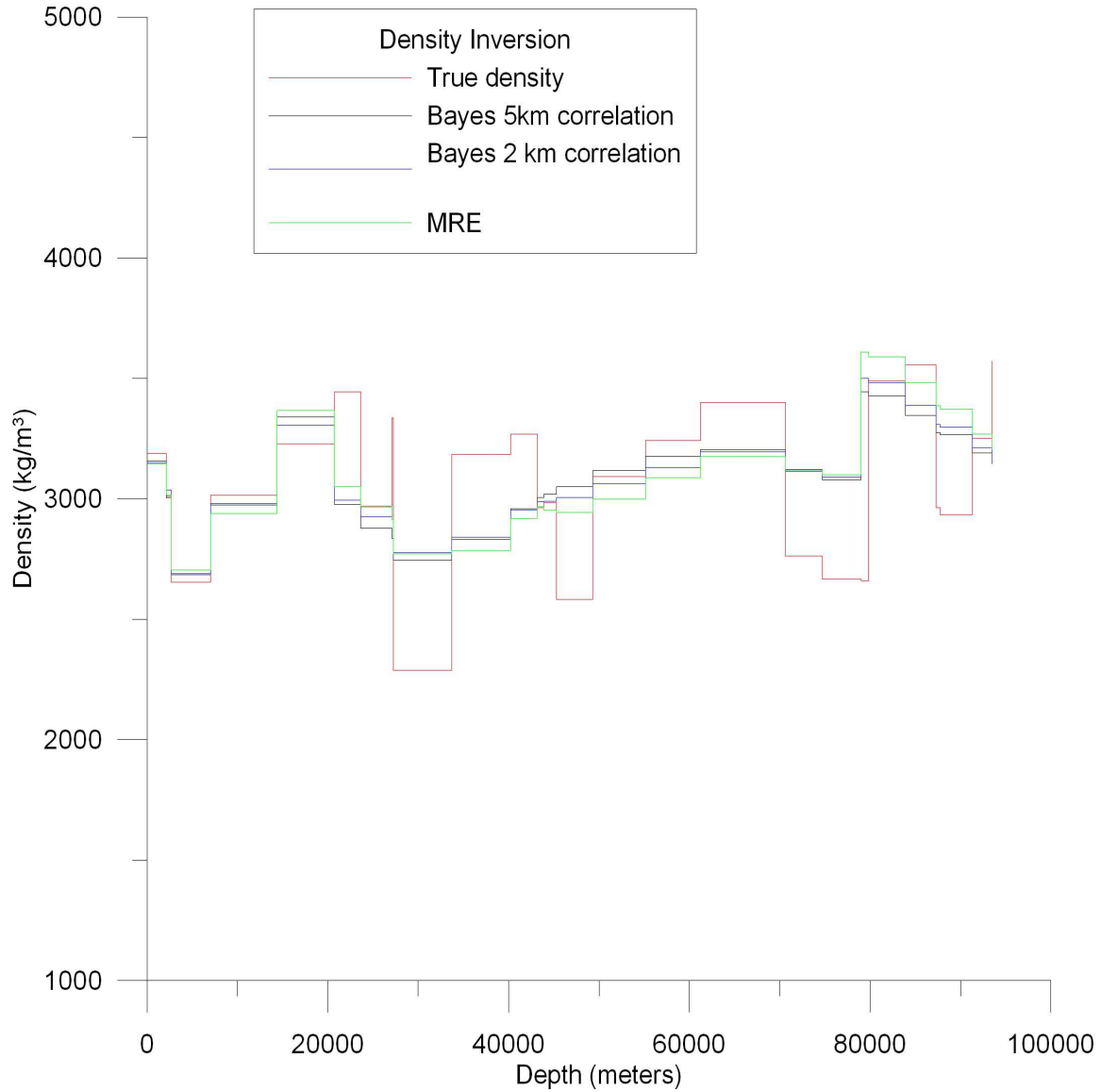


Figure 1. Comparison of linear inverse results over a layered fault.

461 equated to the observed data, and we only enforce a condition of known covariance \mathbf{C}_q equal
 462 to what would be obtained from Bayes. If the unknown matrix of Lagrange multipliers is set
 463 equal to the observed noise error \mathbf{C}_d^{-1} then we duplicate the Bayes solution. It is important
 464 to note though that this argument is circular.

465 In general though, since the input data are known only to within a certain accuracy, it
 466 is important that any inversion method allow for errors in the measured data. The MRE
 467 approach can accommodate such uncertainty and this subject has been discussed by Johnson
 468 and Shore (1984) and Ulrych et al. (1990). We show in this paper how a Gaussian prior can be

469 updated with new data such that $\bar{\mathbf{d}}$ does not equal (exactly) the observed data. The important
470 fact from the point of view of MRE is that the form of the solution with uncertain data is not
471 changed; only a modification of the standard solution technique is required.

472 The classic results of the posterior solution under Gaussian priors and likelihood is re-
473 peated for clarity and completeness. A variety of MRE solutions are reproduced under a
474 number of assumed moments and these stop at second-order central moments. Various so-
475 lutions of Jacobs and van der Geest (1991) were repeated and clarified. Menke's weighted
476 minimum length solution was shown to have a basis in information theory, and the classic
477 least squares estimate is shown as a solution to MRE under the conditions of more data than
478 unknowns and where we utilize the observed data and their associated noise.

479 One of the first questions posed in this paper was under what circumstances should one use
480 entropic rather than Bayesian methods? The answer to that question may not be settled here.
481 Certainly maximum entropy and MRE have had a huge success in geophysics (see Shore, 1981;
482 Ulrych and Bishop, 1975). The MRE approach in the author's opinion is very flexible and may
483 be much less demanding than the full Bayesian solution. MRE has shown to be ideal when
484 constraining moments are known, for example: known second moments from autocorrelation
485 (Shore, 1981), gross earth properties and density inversion (Woodbury and Ulrych, 1998b),
486 and hydrologic applications (Kaplan et al., 2002; Singh, 2000). A very good case can be made
487 for MRE in cases where little or no prior information is available (Kennedy et al., 2000)
488 because MRE enforces conditions of independence. When some measure on the error on a
489 set of observations is known the MRE approach can accommodate such uncertainty. Having
490 said that, so can empirical Bayes, when say, the noise variance or other hyperparameters are
491 unknown and have to be estimated from the data themselves.

492 If the prior in inverse problems has an important role then one could ask what are the
493 ramifications associated with a particular choice? Certainly it is somewhat unpleasant to
494 have to introduce some prior information to a problem in the first place to obtain a solution.
495 As Menke (1989) suggested, the importance of the prior information depends on the use
496 one plans for the results of the inversion. If one simply wants an exploration target then
497 perhaps that choice is not important. If one plans on using the results of the inversion,
498 specifically the errors in the estimates, then the validity of the prior assumptions (and pdfs
499 in a probabilistic inversion) are critically important. This may be the case when the inverse
500 problems are conducted in the realm of the environmental sciences or in engineering risk
501 assessment. In these cases MRE may indeed play a important role in inverse problems in
502 which it is desired to be maximally uncommitted with respect to unknown information.

503 It is important to note that (Seth and Kapur, 1990)

504 “For the same problem, different methods of estimation can lead, quite expectedly to different results.
505 Since statistical inference is inductive, there can be no perfect solution and there can be even differences
506 of opinion as to which one is best.”

507 **6 ACKNOWLEDGMENTS**

508 This article is meant as a tribute to Professor Jagat Kapur who passed away in September
509 4, 2002 at the age of 78 in Delhi. He was a pioneer in the field of maximum entropy who
510 sought after a deep understanding of probabilities and the commonality in statistical inference
511 approaches. He was giant in the field and he left an outstanding legacy of teaching, research
512 and mentorship.

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